

A 2D-Non-Parabolic Six Moments Model

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The most accurate way to describe carrier transport is to solve the Boltzmann transport equation (BTE), for instance with the very time consuming Monte-Carlo (MC) technique. On an engineering level however macroscopic transport models are more efficient. Multiplication of the BTE with weight functions, approximation of the scattering integral with a macroscopic relaxation time and integration over k-space yields, for instance, the drift-diffusion model, the energy transport model, and the six moments model [1]. As demonstrated in [2] and [3] for the bulk case, a non-parabolic six moments model may be accurate down to 30 nm. The challenge is to model higher-order transport parameters like the energy relaxation time τ_1 , the second-order relaxation time τ_2 , the energy mobility μ_1 and the second-order mobility μ_2 (see Figure 2) with as few simplifying assumptions as possible. A good choice is the use of tabulated data extracted from MC simulations [4]. So far, bulk data has been used to examine higher-order parameters in a device. However, important effects like surface roughness scattering or the quantization in inversion layers are not included in bulk MC-data.

In order to account for these effects, a selfconsistently coupled Subband-Monte-Carlo simulator (SMC) and a Schrödinger-Poisson (SP) solver (see Figure 1) are applied. The SP solves the quantum confinement and the SMC simulates the 2D transport in each subband [5]. In order to be consistent with the 2D-parameters the approach consistently, we have developed a 2D-higher-order-transport model. The following equations yield the 2D and 3D transport model with D as the dimension of the system. With $D_0 = 1$, $D_1 = 1$, $D_2 = D/2$, $D_3 = (2 + D)/2$, $D_4 = (2 + D)D/2^{4-D}$, and $D_5 = (4 + D)(2 + D)/4$ the model reads

$$\begin{aligned}
 \phi_0 : D_0 \partial_t n - \frac{1}{q_0} \nabla_r \mathbf{J}_n &= 0, \\
 \phi_2 : D_2 k_B \partial_t (n T_n) + \nabla_r \mathbf{S}_n - \mathbf{E} \cdot \mathbf{J}_n &= -D_2 n k_B \frac{T - T_0}{\tau_1}, \\
 \phi_4 : D_4 \frac{k_B^2}{m^*} \partial_t (n T_n^2 \beta) + \frac{2}{m^*} \nabla_r \mathbf{K}_n + \frac{4q_0}{m^*} \mathbf{E} \cdot \mathbf{S}_n &= -D_4 \frac{n}{m^*} k_B^2 \frac{T^2 - T_0^2}{\tau_2}, \\
 \phi_1 : \mathbf{J}_n &= \mu_0 H_0 (D_1 \nabla_r (k_B n T_n) + h_0 q_0 \mathbf{E} n), \\
 \phi_3 : \mathbf{S}_n &= -\mu_1 H_1 \left(D_3 \nabla_r \left(n \frac{(k_B T_n)^2}{q} \beta \right) + h_1 D_3 k_B T_n \mathbf{E} n \right), \\
 \phi_5 : \mathbf{K}_n &= -\mu_2 H_2 \left(D_5 \nabla_r \left(n \frac{(k_B T_n)^3}{q} \gamma \right) + h_2 D_5 (k_B T_n)^2 \mathbf{E} n \beta \right).
 \end{aligned}$$

Here, ϕ denotes the moment of the equation. The even moments are the balance equations and the odd ones are the fluxes. H_i are the non-parabolicity factors. For parabolic bands, $H_i = 1$. β is the kurtosis and denotes the deviation from the Maxwellian distribution function (see Figure 4). For the 2D and the 3D case the kurtosis is defined as:

$$\beta = \frac{D}{D+2} n \frac{\langle \epsilon^2 \rangle}{\langle \epsilon \rangle^2}$$

In Figure 3 we show the output characteristics of the 2D drift-diffusion, energy transport, and the six moments model of a SOI-MOSFET.

We introduce a 2D-nonparabolic-six-moments model for scaled devices. Due to the parameter modeling based on MC-tables, important effects like quantization as well as surface roughness scattering are inherently included.

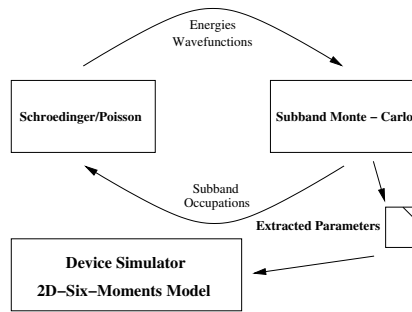


Figure 1: Transport parameters of a 2D-electron gas in an inversion layer are extracted selfconsistently and modeled through a whole device with a device simulator.

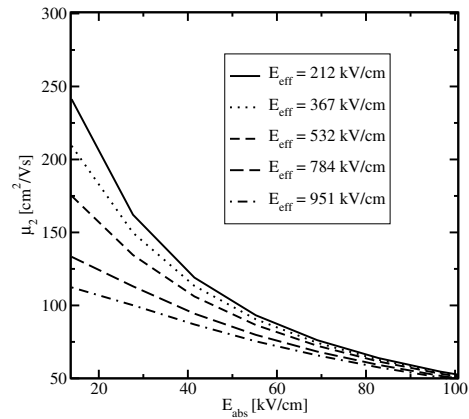


Figure 2: Extracted second-order mobility as a function of the driving field for different effective fields. The device simulator interpolates between these curves.

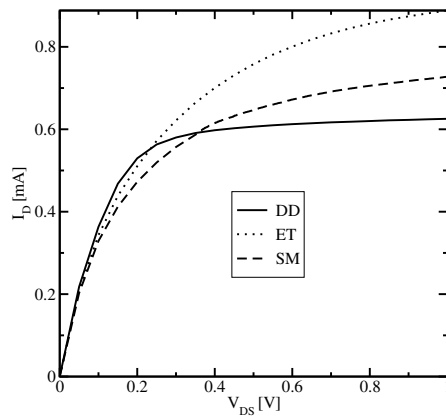


Figure 3: The output characteristic of a UTB SOI-Mosfet with a gate length of 40 nm calculated with the drift-diffusion model, the energy transport model, and the six moments model.

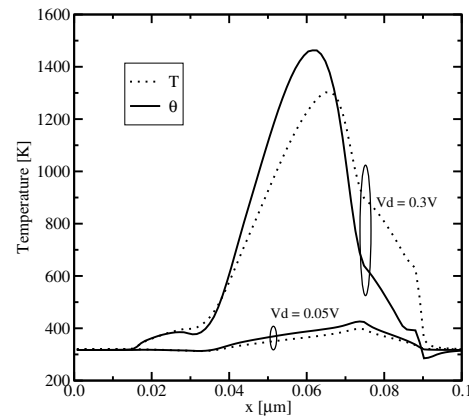


Figure 4: The second-order temperature $\theta = \beta T$ in comparison to the carrier temperature T . With increasing drain voltage, the deviation from the Maxwellian distribution ($\theta \approx T$, $\beta \approx 1$) increases.

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