## A Rigorous Study of Nanoscaled Transistors Based on Single-Layer MoS<sub>2</sub>

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Since the successful experimental isolation of graphene in 2004, ultra-thin two-dimensional structures are being widely studied as potential building blocks for future electronic devices. Among various two-dimensional materials, single-layer (SL) MoS<sub>2</sub> has attracted much attention. For a SL of MoS2 a direct bandgap of 1.8 to 1.9 eV has been reported [1], which is suitable for various electronic applications. Recently, FETs based on SL MoS<sub>2</sub> with an  $I_{on}/I_{off}$ ratio as high as  $\sim 10^8$  and a sub-threshold swing of  $\sim 70$  mV/decade have been achieved [2]. However, the reported mobility is below that of ultra thin body or strained Si and III-V materials. It is believed that extrinsic sources such as charged impurities (CI) [3] and inevitable Schottky contacts [4] limit the characteristics of devices based on SL MoS<sub>2</sub>. High- $\kappa$  gate insulators, such as HfO<sub>2</sub>, can reduce CI scattering effects and boost the mobility [3, 5], however, they can degrade the mobility due to remote phonon (RP) scattering [6]. The source of this scattering is in the surrounding dielectrics via long-range Coulomb interactions, provided that the dielectrics support polar vibrational modes.

To study electronic transport in SL MoS<sub>2</sub> we solved the NEGF equations self-consistently with the Poisson equation based on the box integration method. An effective mass of  $m^* = 0.48m_0$  has been assumed for both longitudinal and transverse directions [7]. We have considered intrinsic electron-phonon interactions including the longitudinal acoustic (LA), the transverse acoustic (TA), the longitudinal optical (LO), and polar optical phonons (POP) with the parameters adopted from Ref. [7]. The mobility is calculated based on the method explained in Refs. [8, 9]. For device simulation we assumed a channel length of 20 nm, a 10 nm thick HfO<sub>2</sub> gate insulator, and a 50 nm thick

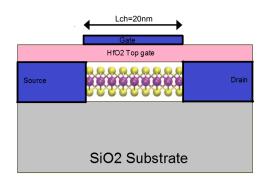


Fig. 1. Schematic view of the simulated device. An n-type top gate SL MoS<sub>2</sub> FET with ohmic contacts. The channel length is 20 nm and the gate oxide is 10 nm thick HfO<sub>2</sub> layer with  $\kappa=22$ . The substrate is assumed to be SiO<sub>2</sub>.

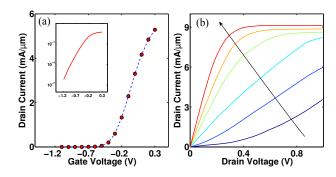


Fig. 2. (a) The transfer characteristics in the presence of scattering at  $V_{\rm DS}$ =0.1 V. Transconductance is  $g_m=9.6~{\rm mS}/\mu$  m. The inset shows transfer characteristics in logarithmic scale. (b) The output characteristics for  $V_{\rm GS}$ = - 0.6 to +0.4 V with 0.2 V step (the arrow indicates the direction of  $V_{\rm GS}$  increase). Current saturation is observed for  $V_{\rm DS}>0.3~{\rm V}$ .

 ${
m SiO_2}$  substrate, see Fig. 1. In this study the dynamic screening of remote phonon modes stemmed from HfO<sub>2</sub> and static screening of charged impurities are included. Figure 2 shows the transfer and output characteristics in the presence of intrinsic phonon scattering. The results indicate a relatively high  $I_{\rm on}$ 

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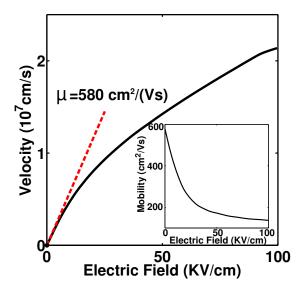


Fig. 3. The velocity as a function of the applied electric field. The inset shows the mobility as a function of the applied electric field. The low field mobility is about  $580 \text{ cm}^2/\text{Vs}$  and decreases to  $130 \text{ cm}^2/\text{Vs}$ ) under high electric fields.

of 9 mA/ $\mu$ m and a high  $I_{\rm on}/I_{\rm off}$  ratio of about  $10^7$ , which are close to the ballistic limit. Low field phonon-limited mobility is evaluated to be  $580~{\rm cm^2/(Vs)}$  in good agreement with the result of Ref. [7]. The effect of high fields on the mobility and carrier velocity are depicted in Fig. 3. The significant drop in the mobility is due to increased polar and non-polar optical phonon scattering at high electric fields.

Figure 4 shows the effect of CI scattering and RP scattering for an average carrier density of  $\sim 10^{13} \text{ cm}^{-2}$ . By using a high- $\kappa$  insulator the CI scattering is suppressed which enhances the mobility. On the other hand, high- $\kappa$  insulators introduce RP scattering which in turn reduce the mobility. Therefore, at low CI concentrations the mobility for a device with SiO<sub>2</sub> as the top gate is higher than that with HfO<sub>2</sub>. Figure 4 shows the evaluated mobility in the presence of intrinsic phonons (IP), CI, and RP (CI+RP) is in good agreement with experimental results from Refs. [3, 10]. Table I compares the mean free path and the mobility for each scattering mechanism studied in this work. The results show that acoustic phonons play a more significant role in short channel devices than other intrinsic phonon modes. However, RP scattering due to a 30 nm high- $\kappa$  HfO<sub>2</sub> results in

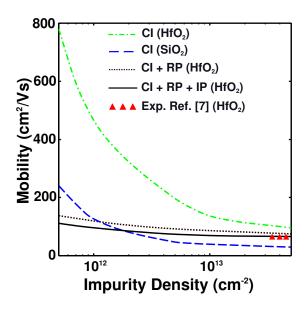


Fig. 4. Charged impurity (CI) and remote phonon (RP) limited mobility as a function of charged impurity density for 30nm thick HfO<sub>2</sub> and also SiO<sub>2</sub> top gate insulator. The mobility due to IPs+CI+RP with HfO<sub>2</sub> gate insulator is in good agreement with experimental results in Refs. [3, 10].

the smallest mean free path. The presented results can be used for appropriate selection of the gate insulator material for optimal device performance.

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TABLE I  $\label{thm:lower}$  The mean free path ( $\lambda$ ) and low-field mobility ( $\mu$ ) for each scattering source. RP scattering is calculated for for a 30 nm thick HfO $_2$ .

Phonon Modes	λ (nm)	$\mu \text{ (cm}^2/\text{(Vs))}$
LA+TA	35	923
LO	157	4161
POP	87	2316
RP	5.7	151

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